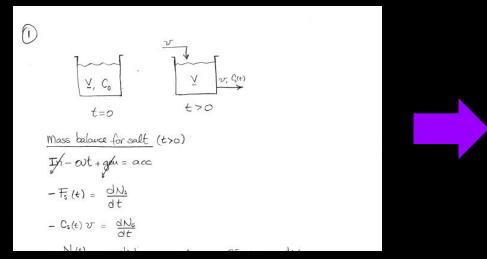
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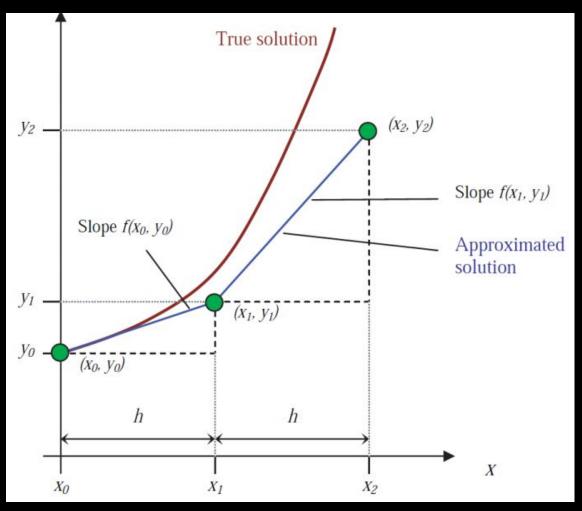




ChE352 Numerical Techniques for Chemical Engineers Professor Stevenson

Lecture 10

Recall: Initial Value Problems



$$\frac{dy}{dt} = f(t, y)$$

$$a \le t \le b$$
, $y(a) = \alpha$

$$t, y \in \mathbb{R}, \quad y : \mathbb{R} \to \mathbb{R}$$

$$f: \mathbb{R}^2 \to \mathbb{R}$$

f continuous

Why can't we use trapezoidal integration? What method can we use instead?

Can you find more IVP examples?

- Anything involving <u>rate of change</u>
 - Reaction rates
 - F = ma
 - Epidemics
 - Time-dependent Schrodinger equation

$$i\hbarrac{\partial}{\partial t}\Psi(x,t)=\left[-rac{\hbar^2}{2m}rac{\partial^2}{\partial x^2}+V(x,t)
ight]\Psi(x,t)$$

- Other examples?
- We define dy/dt = f(t, y) because f(t, y) is the function we actually <u>have</u> in IVPs
 - y is the function we want

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Recall: Euler's Method

$$y(t_{i+1}) \approx y(t_i) + hf(t_i, y(t_i))$$

Pronounced the same as "oiler"
Solve the IVP by taking steps along the derivative

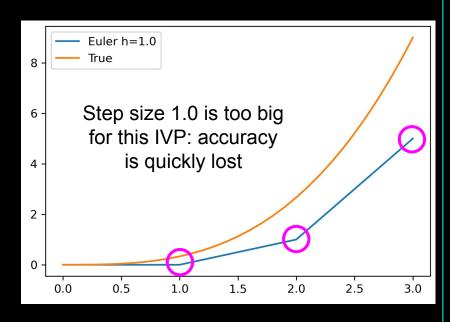


Recall: Euler's Method

$$y(t_{i+1}) \approx y(t_i) + hf(t_i, y(t_i))$$

Example:
$$f(t) = t^2$$
 $t_0 = 0$ $y(t_0) = 0$ $h = 1$
 $t_1 = h = 1$
 $y(1) \approx 0 + 1 * 0^2 = 0$
 $t_2 = 2 * h = 2$
 $y(2) \approx 0 + 1 * 1^2 = 1$
 $t_3 = 3 * h = 3$
 $y(3) \approx 1 + 1 * 2^2 = 5$

Step size 1.0 is too big for this IVP: accuracy is quickly lost



Euler's Method to get all w[i]

We can define a vector of "time" (call it "t") and calculate our approximate y(t) (aka "w") by iterating forwards in "time" from $t_0 = a$:

$$t = \begin{bmatrix} t_0 \\ t_1 \\ t_2 \\ \vdots \\ t_{N-1} \end{bmatrix} = \begin{bmatrix} a \\ a+h \\ a+2h \\ \vdots \\ a+(N-1)h \end{bmatrix}, \quad w = \begin{bmatrix} w_0 \\ w_1 \\ w_2 \\ \vdots \\ w_{N-1} \end{bmatrix} = \begin{bmatrix} y(t_0) \\ w_0 + hf(t_0, w_0) \\ w_1 + hf(t_1, w_1) \\ \vdots \\ w_{N-2} + hf(t_{N-2}, w_{N-2}) \end{bmatrix}$$

 $t, w \in \mathbb{R}^N$ Why "time" in quotes?

What if we want in-between w values?

Activity: Euler in Python (15 minutes)

$$y(t_{i+1}) \approx y(t_i) + hf(t_i, y(t_i))$$
 Euler's method

Write a Python function which implements Euler's method for the IVP for this reaction:

$$\frac{dC_{EB}}{d\tau} = -k_f C_{EB}, \quad C_{EB} \left(\tau = 0\right) = C_{EB}^o$$

Assume: $k_f = 1.0$, $C^0_{EB} = 2.0$, $\tau_{final} = 10.0$

Use step size h = 0.01. Does the h value matter?

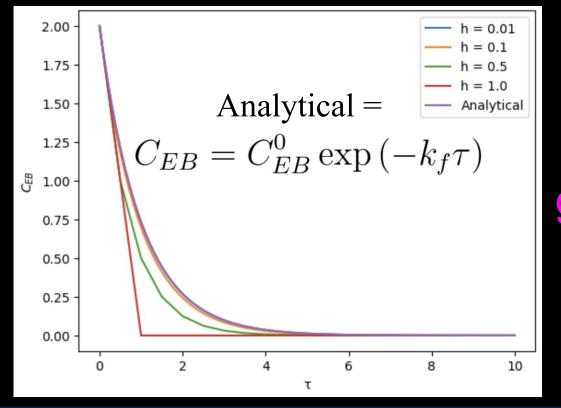
Make a list of your approximate C_{EB} at each step, and if you have time, plot your results vs t

Solution: Euler in Python

$$y(t_{i+1}) \approx y(t_i) + hf(t_i, y(t_i))$$
 Euler's method

$$rac{dC_{EB}}{d au} = -k_f C_{EB}, \quad C_{EB} \left(au = 0\right) = C_{EB}^o$$

Euler solution is nearly exact at small dt



Euler solution goes bad fast at large dt

Is there a better IVP method?

- Euler's Method is straightfoward, works if you can afford a small h
 - Local error O(h²), global error O(h)
- But we want better than O(h)
- What is local error vs global error?
- Why is global error 1/h times bigger?
- Why can't we always make h smaller?
- How can we make a better method?
 - Consider where Euler's Method comes from

Taylor Methods of Order n

 Euler's method uses just the linear Taylor terms, but we could use up to any n:

$$y(t) = y(t_{i}) + (t - t_{i})y'(t_{i}) + \frac{(t - t_{i})^{2}}{2}y''(t_{i}) + \dots + \frac{(t - t_{i})^{n}}{n!}y^{(n)}(t_{i}) + \frac{(t - t_{i})^{n+1}}{(n+1)!}y^{(n+1)}(\xi_{i})$$

$$\Rightarrow y_{i+1} = y_{i} + hf(t_{i}, y_{i}) + \frac{h^{2}}{2}f'(t_{i}, y_{i}) + \dots + \frac{h^{n}}{n!}f^{(n-1)}(t_{i}, y_{i}) + \frac{h^{n+1}}{(n+1)!}f^{(n)}(\xi_{i}, y(\xi_{i}))$$
Linear terms Quadratic and higher terms Error term

- By definition, y'(t) = f(t, y) We always have this in an IVP
- 2nd derivative: y''(t) = f'(t, y) Might not have this
- n-th derivative: $y^n(t) = f^{n-1}(t, y)$ Good luck

Taylor Methods of Order n

 Euler's method uses just the linear Taylor terms, but we could use up to any n:

$$y(t) = y(t_{i}) + (t - t_{i})y'(t_{i}) + \frac{(t - t_{i})^{2}}{2}y''(t_{i}) + \dots + \frac{(t - t_{i})^{n}}{n!}y^{(n)}(t_{i}) + \frac{(t - t_{i})^{n+1}}{(n+1)!}y^{(n+1)}(\xi_{i})$$

$$\Rightarrow y_{i+1} = y_{i} + hf(t_{i}, y_{i}) + \frac{h^{2}}{2}f'(t_{i}, y_{i}) + \dots + \frac{h^{n}}{n!}f^{(n-1)}(t_{i}, y_{i}) + \frac{h^{n+1}}{(n+1)!}f^{(n)}(\xi_{i}, y(\xi_{i}))$$

$$y_{i+1} \approx y_i + hf(t_i, y_i) + \frac{h^2}{2} f'(t_i, y_i) + \dots + \frac{h^n}{n!} f^{(n-1)}(t_i, y_i)$$

- If we use a series of order n, the <u>local error</u> for each step is $O(h^{n+1})$ (Why?)
- Global error after all steps is $O(h^n)$ (Why?)

Activity: 2nd Order Taylor Methods

$$y_{i+1} \approx y_i + hf(t_i, y_i) + \frac{h^2}{2}f'(t_i, y_i) + \dots + \frac{h^n}{n!}f^{(n-1)}(t_i, y_i)$$

Translate the Taylor polynomial formula above into an iterative step for the 2^{nd} -order Taylor method for IVPs, giving w_{i+1} in terms of w_i , t_i , f, f', and h.

Use your general expression to define the iterative step $w_{{}_{i+1}}$ for this IVP:

$$y' = y - t$$
 $t_0 = 0$ $y(0) = e + 1$

Leave your expression in terms of h (Why?)

Answer: 2nd Order Taylor Methods

Euler's method:

$$y' = y - t$$
 $t_0 = 0$ $y(0) = e + 1$

$$w_0 = e + 1$$

 $w_{i+1} = w_i + h(w_i - t_i) = (h+1)w_i - ht_i \quad (i = 0...N-2)$

2nd order Taylor:

$$w_{i+1} = w_i + hf(t_i, w_i) + \frac{h^2}{2}f'(t_i, w_i) = w_i + h(w_i - t_i) + \frac{h^2}{2}\frac{d}{dt}(w - t)_i$$

$$= w_i + hw_i - ht_i + \frac{h^2}{2}(w_i - t_i) - \frac{h^2}{2} = \left[\frac{h^2}{2} + h + 1\right]w_i - h\left(\frac{h}{2} + 1\right)t_i - \frac{h^2}{2}$$

What are some drawbacks of this method?

The problem with f'

- Taylor methods gain more accuracy by using more derivatives of f
 - \circ Recall: $y^{n}(t) = f^{n-1}(t, y)$
- But derivatives of f are rarely available
- Can we <u>approximate</u> f'(t, y) using the values of f(t, y)? How?
- The resulting methods are the most popular IVP solvers: Runge-Kutta

Runge-Kutta Methods: RK2

Use 2D Taylor series & the chain rule to **find** $f'(t_i, y_i)$, with $\Delta t = h/2$ and $\Delta y = \Delta t f(t_i, y_i)$. Then plug $f'(t_i, y_i)$ into the 2^{nd} order Taylor method.

$$f\left(t + \Delta t, y + \Delta y\right) \approx f\left(t, y\right) + \left[\Delta t \left(\frac{\partial f}{\partial t}\right)_{t, y} + \Delta y \left(\frac{\partial f}{\partial y}\right)_{t, y}\right]$$
2D Taylor series in y, t

Chain rule gives
$$f'(t_i, y_i)$$
 = $\left(\frac{\partial f}{\partial t}\right)_{t_i, y_i}$ + $\left(\frac{\partial f}{\partial y}\right)_{t_i, y_i} \left(\frac{dy}{dt}\right)_{t_i}$

$$y_{i+1} \approx y_i + hf(t_i, y_i) + \frac{h^2}{2} \boxed{f'(t_i, y_i)} \begin{array}{l} \text{2nd order} \\ \text{Taylor method} \\ \text{needs } f'(t_i, y_i) \end{array}$$

Runge-Kutta Methods: RK2

$$f(t_{i+1}, y_{i+1}) \approx f(t_i, y_i) + \Delta t \left(\frac{\partial f}{\partial t}\right)_{t_i, y_i} + \Delta y \left(\frac{\partial f}{\partial y}\right)_{t_i, y_i}$$
 Same as chain rule!
$$= f(t_i, y_i) + \frac{h}{2} \left[\left(\frac{\partial f}{\partial t}\right)_{t_i, y_i} + \frac{h}{2} f(t_i, y_i)\right] \left(\frac{\partial f}{\partial y}\right)_{t_i, y_i}$$
 Same as chain rule!
$$= f(t_i, y_i) + \frac{h}{2} \left[\left(\frac{\partial f}{\partial t}\right)_{t_i, y_i} + \frac{f(t_i, y_i)}{\partial y}\right] \left(\frac{\partial f}{\partial y}\right)_{t_i, y_i} = f(t_i, y_i) + \frac{h}{2} \left[\left(\frac{\partial f}{\partial t}\right)_{t_i, y_i} + \frac{dy}{dt}\right] \left(\frac{\partial f}{\partial y}\right)_{t_i, y_i}$$

$$= f'(t_i, y_i) \approx \frac{2}{h} \left[f(t_{i+1}, y_{i+1}) - f(t_i, y_i)\right]$$

Runge-Kutta Methods: RK2

$$f'(t_{i}, y_{i}) \approx \frac{2}{h} \Big[f(t_{i+1}, y_{i+1}) - f(t_{i}, y_{i}) \Big], \qquad \text{Given } f', \text{ we can plug} \\ w_{i+1} = w_{i} + hf(t_{i}, w_{i}) + \frac{h^{2}}{2} f'(t_{i}, w_{i}) \qquad \text{Taylor IVP method} \\ w_{i+1} = w_{i} + hf(t_{i}, w_{i}) + \frac{h^{2}}{2} \Big(\frac{2}{h} \Big) \Big[f(t_{i+1}, w_{i+1}) - f(t_{i}, w_{i}) \Big] \\ = w_{i} + hf(t_{i}, w_{i}) + h \Big[f(t_{i+1}, w_{i+1}) - f(t_{i}, w_{i}) \Big] \rightarrow \\ w_{i+1} = w_{i} + hf(t_{i}, w_{i}) + hf(t_{i+1}, w_{i+1}) - hf(t_{i}, w_{i}) \\ = w_{i} + hf(t_{i+1}, w_{i+1}) = w_{i} + hf(t_{i+1}, w_{i+1}) \rightarrow f(t_{i}, w_{i}) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) = w_{i} + hf(t_{i}, w_{i}) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) = w_{i} + hf(t_{i}, w_{i}) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) = w_{i} + hf(t_{i}, w_{i}) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) = w_{i} + hf(t_{i}, w_{i}) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1})) \rightarrow \\ (1 + hf(t_{i+1}, w_{i+1}$$

$$\left| w_{i+1} = w_i + hf\left(t_i + \frac{h}{2}, w_i + \frac{h}{2}f\left(t_i, w_i\right)\right) \right| \leftarrow \frac{\text{RK2, aka "midpoint method for IVPs"}}{\text{method for IVPs"}}$$

Activity: RK2 in Python (10 minutes)

$$y(t_{i+1}) \approx y(t_i) + hf(t_i, y(t_i))$$
 Euler's method

Copy your Python IVP solver from before and change it to RK2:

$$w_{i+1} = w_i + hf\left(t_i + \frac{h}{2}, w_i + \frac{h}{2}f(t_i, w_i)\right)$$
 RK2

Make a list of your approximate C_{EB} at each step, and if you have time, plot your results vs t

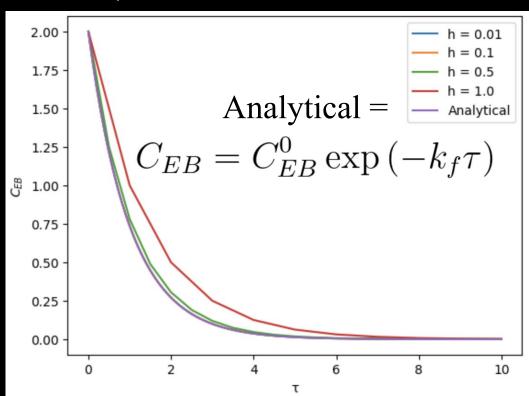
How does the dependence on h change?

Solution: RK2 in Python

$$w_{i+1} = w_i + hf\left(t_i + \frac{h}{2}, w_i + \frac{h}{2}f(t_i, w_i)\right)$$
 RK2

$$\frac{dC_{EB}}{d au} = -k_f C_{EB}, \quad C_{EB} \left(au = 0 \right) = C_{EB}^o$$

RK2 solution is nearly exact at small dt



RK2 solution does not go bad so fast

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Better Runge-Kutta?

- Different values for Δt and Δy in 2D Taylor make new IVP methods (F&B 185-187)
- Order 2 methods have global approximation error of O(h²)
- Most common RK method for solving IVPs is order 4, which uses the Taylor terms up to h⁴
- This method is called <u>RK4</u> or just <u>The</u> <u>Runge-Kutta Method</u> for IVPs
- Given this description, what is the big-O of local & global error for RK4?

"The" Runge-Kutta Method: RK4

$$w_{i+1} = w_i + \frac{1}{6} (k_1 + 2k_2 + 2k_3 + k_4)$$

Where:
$$k_1 = hf(t_i, w_i)$$

$$k_2 = hf\left(t_i + \frac{h}{2}, w_i + \frac{1}{2}k_1\right)$$

$$k_3 = hf\left(t_i + \frac{h}{2}, w_i + \frac{1}{2}k_2\right)$$

$$k_4 = hf\left(t_{i+1}, w_i + k_3\right)$$

- Like RK2 but more
- Global error O(h⁴)
- Requires 4 calls to f(t, y) per step
- Don't need f'(t, y)
- Usually the sweet spot for accuracy

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Why stop at RK4?

- The main cost for using an IVP algorithm is the calls to function f – fewer is better
- Euler needs 1 function evaluation per step
- RK4 needs 4
- RK4 is only useful if it allows step sizes over 4x bigger, with the same accuracy (it does)
- Table on p. 188 of F&B shows that <u>RK4 is</u> <u>superior to lower and higher order methods</u> by this metric under reasonable assumptions

Activity: Local Error in RK4

1. Use RK4 to estimate y(0.1) for this IVP:

$$y' = y - t$$
 $t_0 = 0$ $y(0) = e + 1$ $h = 0.1$

2. Just as a demonstration of the error, compare your approximation to the exact answer y(t) = e^{t+1} + t + 1 to get the actual local relative approximation error. Is it similar in scale to h⁵?

Answer: Local Error in RK4

$$w_0 = e + 1, \quad h = 0.1$$

$$w_1 = e + 1 + \frac{1}{6} (k_1 + 2k_2 + 2k_3 + k_4) = \boxed{4.104165794}$$

Where:
$$k_1 = 0.1[w_0 - t_0] = 0.1e + 0.1$$

$$k_2 = (0.1) \left[w_0 + \frac{1}{2} k_1 - t_0 - \frac{h}{2} \right] = 0.105e + 0.1$$

$$k_3 = (0.1) \left[w_0 + \frac{1}{2} k_2 - t_0 - \frac{h}{2} \right] = 0.10525e + 0.1$$

$$k_4 = (0.1)[w_0 + k_3 - t_1] = 0.110525e + 0.1$$

$$y_1 = e^{1.1} + 1.1 \rightarrow error = 5.61 \times 10^{-8}$$
 (really tiny)

SciPy generic IVP solver: solve_ivp

```
from scipy.integrate import solve_ivp
sol = solve_ivp(fun, (t0, t_end), [y0])
plt.plot(sol.t, sol.y[0], label='RK45')
```

- Uses RK4 but with dynamic h, with an error estimate based on RK5 - known as RK4(5)
 - Also has other, specialized methods
- Can solve for multi-dimensional y in f(t, y)
- Returns an object containing data about the solution, including sol.t, sol.y, & sol.success

IVP Systems

 1D problems are common, but so are IVPs with multiple outputs:

$$\frac{dN_{S}}{dz} = \frac{k_{f}N_{EB}}{v} - \frac{k_{r}N_{S}N_{H}}{v} = R_{S}', \quad N_{S}(z=0) = vC_{S}'$$
Where are the dependent $\frac{dN_{EB}}{dz} = -R_{S}', \quad N_{EB}(z=0) = vC_{EB}'$

$$\frac{dP}{dz} = -\frac{\rho v^{2}}{d_{p}} \left(\frac{1-\varepsilon}{\varepsilon^{3}}\right) \left[\frac{150(1-\varepsilon)}{\text{Re}_{p}} + 1.75\right], \quad \frac{dN_{W}}{dz} = 0, \quad N_{S} = N_{H}$$

$$P\hat{V} = ZRT \implies \rho = \frac{PM}{ZRT} \implies v = \frac{ZRT}{P}(N_{EB} + N_{S} + N_{H} + N_{W})$$

 We need output to be a <u>vector</u> instead of a scalar - u now instead of y

Numerical Soln. of IVP Systems

Suppose your problem now looks like this:

$$\frac{du_1}{dt} = f_1(t, u_1, u_2, \dots, u_m), \quad u_1(t = t_0) = a_1$$

$$\frac{du_2}{dt} = f_2(t, u_1, u_2, \dots, u_m), \quad u_2(t = t_0) = a_2$$

$$\vdots$$

$$\frac{du_m}{dt} = f_m(t, u_1, u_2, \dots, u_m), \quad u_m(t = t_0) = a_m$$

$$\frac{du(t)}{dt} = f(t, u(t)),$$

$$u(t = t_0) = a,$$

$$t_0 \le t \le t_{\text{max}}$$

$$u : \mathbb{R} \to \mathbb{R}^m,$$

$$f : \mathbb{R}^{m+1} \to \mathbb{R}^m,$$

$$t \in \mathbb{R}, \quad a \in \mathbb{R}^m$$

Vector function

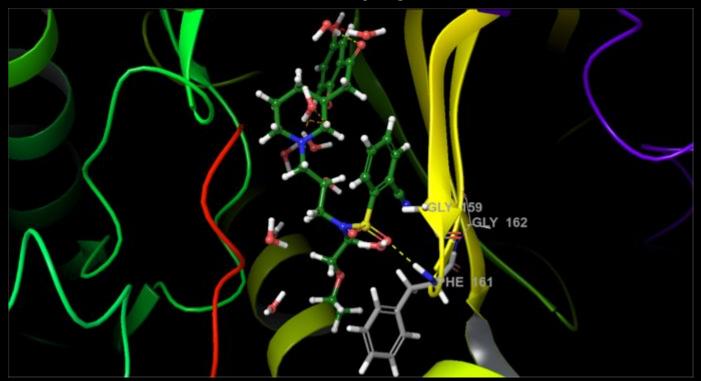
Same methods work!

IVP Systems in Python

```
from scipy.integrate import solve ivp
def fun(t, u): # 3-D IVP
  CA, CB, CC = u
   ... calculate du/dt here ...
   return dAdt, dBdt, dCdt
sol = solve ivp(fun, (t0, t final), u0)
plt.plot(sol.t, sol.y[0], label='[A]')
plt.plot(sol.t, sol.y[1], label='[B]')
plt.plot(sol.t, sol.y[2], label='[C]')
```

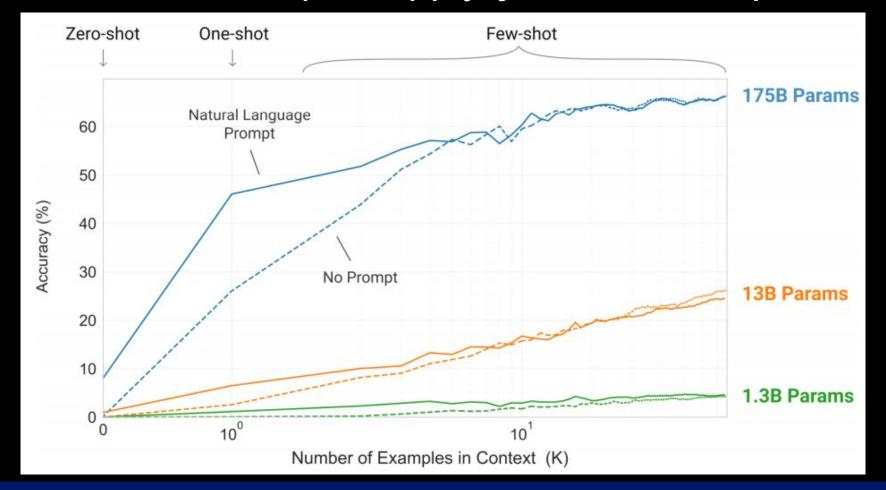
Million+ Dimension IVP Systems

- IVPs often scale to millions of dimensions
- Example: molecular dynamics, every [x, y, z]
 of every atom is another dimension of w(t)
- Same techniques apply, just more compute



10¹²+ Dimension IVP Systems

- Machine learning all known text / images
- Same techniques apply, just more compute



Activity: Coding RK4

• Write a function that $w_0 = y(a) = \alpha$ calculates the next step of RK4: $w_{i+1} = w_i + \frac{1}{6}(k)$

$$w_0 - y(a) - \alpha$$

$$w_{i+1} = w_i + \frac{1}{6} (k_1 + 2k_2 + 2k_3 + k_4)$$

$$k_2 = hf\left(t_i + \frac{h}{2}, w_i + \frac{1}{2}k_1\right)$$

Where: $k_1 = hf(t_i, w_i)$

Try it with this IVP:

def fun(t, w):
$$k_3 = hf\left(t_i + \frac{h}{2}, w_i + \frac{1}{2}k_2\right)$$
 return w - t
t0 = 0; y0 = np.e+1 $k_4 = hf\left(t_{i+1}, w_i + k_3\right)$

When you've got it, compare vs scipy.integrate.solve_ivp

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Pre-reading for next week

Predictor-corrector & adaptive methods for IVPs, higher-order IVPs, stiff IVPs: PNM 22.6-7, F&B 5.6-8.

Verlet integration:

https://www.algorithm-archive.org/contents/verlet_integra tion/verlet_integration.html